

Theoretical possibilities for the existence of tri- and tetraneutron systems

I. V. Simenog^{a,b}, B. E. Grinyuk^a, Yu. M. Bidasyuk^b

^a*Bogolyubov Institute for Theoretical Physics,
Nat. Acad. of Sci. of Ukraine,
Kyiv 03143, Ukraine*

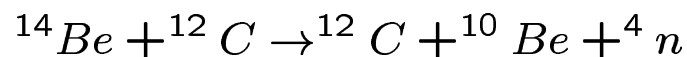
^b*Taras Shevchenko Kyiv National University,
Kyiv 03127, Ukraine*

The main task:

- ▶ To investigate 2-, 3- and 4-neutron systems with different central pairwise n - n potentials and to find such potentials, that can at the same time describe all known two-nucleon data and bind a 4n system or to show that it is impossible.

Motivation:

- ▶ The problem of multineutrons stands for more than 40 years and still there are no final answer
- ▶ Possible observation of free tetraneutron clusters at LPC-Caen in experiments of ${}^{14}\text{Be}$ breakup



(Marqués et al., *Phys. Rev. C*, **65**, 2002)

- ▶ to obtain deeper theoretical knowledge of 3- and 4-fermion systems.

Basic equations (3 fermions)

The total antisymmetric wave function of the 4-fermion system

$$\Psi^a(1, 2, 3) = \frac{1}{\sqrt{2}} (\Phi'(r_1, r_2, r_3)\xi'' - \Phi''(r_1, r_2, r_3)\xi')$$

where Φ' and Φ'' , spatial components with symmetry $[2,1]$ (3n system) are defined as

$$\begin{aligned}\Phi'(r_1, r_2, r_3) &\equiv \Phi(r_1, r_2, r_3), \\ \Phi''(r_1, r_2, r_3) &\equiv \frac{1}{\sqrt{3}} (2\Phi(r_1, r_3, r_2) - \Phi(r_1, r_2, r_3)),\end{aligned}$$

Schrödinger equation for one spatial component of the wave function

$$\left\{ \sum_{i=1}^3 \frac{\mathbf{p}_i^2}{2m} + \frac{1}{2} \sum_{i>j=1}^3 (V_s^+(r_{ij}) + V_t^-(r_{ij})) + \right. \\ \left. + \frac{1}{2} \sum_{(ij) \neq (23)} (-1)^{i+j} (V_s^+(r_{ij}) - V_t^-(r_{ij})) - \right. \\ \left. - \frac{1}{2} \sum_{(ij) \neq (12)} (-1)^{i+j} (V_s^+(r_{ij}) - V_t^-(r_{ij})) \hat{P}_{23} \right\} \Phi = E\Phi$$

Angular momentum of 3-fermion system

In a system of harmonic oscillators we know that the lowest energy state with symmetry [2,1] is the state with angular momentum $L = 1$. Then, to be sure, that we got a lowest state we need to calculate both states with $L = 0$ and $L = 1$ or choose variational function without explicitly set angular momentum.

We use trial variational functions as combination of Gaussians. Variational function for angular momentum $L = 0$:

$$\Phi_{L=0}(\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3) = \hat{A} \sum_{k=1}^N C_k \exp \left(- \sum_{i>j=1}^3 u_{ij}^k r_{ij}^2 \right)$$

and for angular momentum $L = 1$ and odd parity:

$$\Phi_{L=1}(\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3) = \hat{A} \left(q_{1z} \sum_{k=1}^N C_k \exp \left(- \sum_{i>j=1}^3 u_{ij}^k r_{ij}^2 \right) + q_{2z} \sum_{k=1}^M D_k \exp \left(- \sum_{i>j=1}^3 v_{ij}^k r_{ij}^2 \right) \right)$$

q_{1z} and q_{2z} — are z-components of Jacobi relative coordinates.

Basic equations (4 fermions)

The total antisymmetric wave function of the 4-fermion system

$$\Psi^a(1, 2, 3, 4) = \frac{1}{\sqrt{2}} (\Phi'(\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3, \mathbf{r}_4)\xi'' - \Phi''(\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3, \mathbf{r}_4)\xi')$$

where Φ' and Φ'' , spatial components with symmetry [2,2] (4n system) are defined as

$$\begin{aligned}\Phi'(\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3, \mathbf{r}_4) &\equiv \Phi(\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3, \mathbf{r}_4), \\ \Phi''(\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3, \mathbf{r}_4) &\equiv \frac{1}{\sqrt{3}} (2\Phi(\mathbf{r}_1, \mathbf{r}_3, \mathbf{r}_2, \mathbf{r}_4) - \Phi(\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3, \mathbf{r}_4)),\end{aligned}$$

Schrödinger equation for one spatial component of the wave function

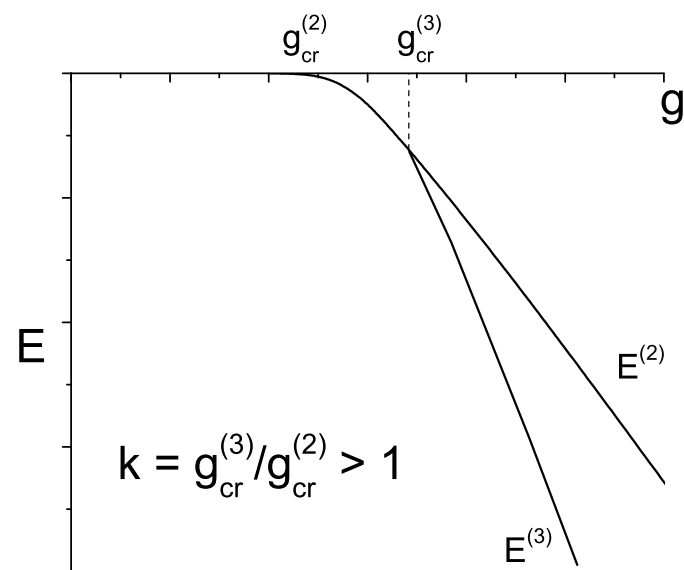
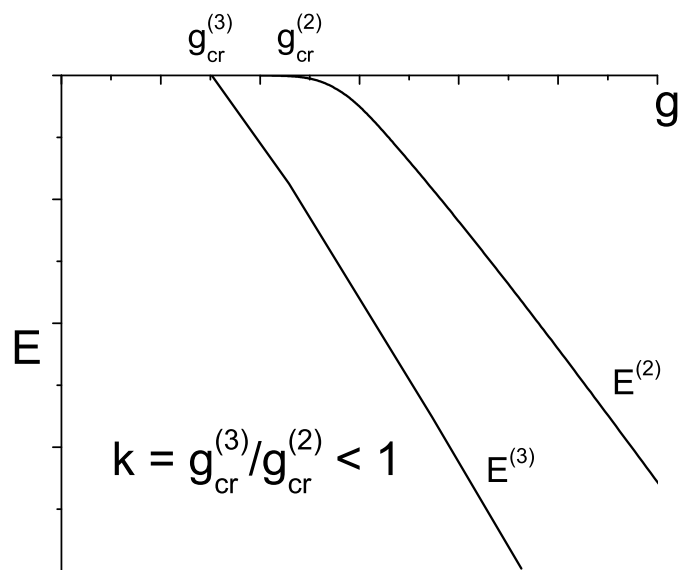
$$\left\{ \sum_{i=1}^4 \frac{\mathbf{p}_i^2}{2m} + \frac{1}{2} \sum_{i>j=1}^4 (V_s^+(r_{ij}) + V_t^-(r_{ij})) + \right. \\ \left. + \frac{1}{2} \sum_{(ij) \neq (14), (23)} (-1)^{i+j} (V_s^+(r_{ij}) - V_t^-(r_{ij})) - \right. \\ \left. - \frac{1}{2} \sum_{(ij) \neq (12), (34)} (-1)^{i+j} (V_s^+(r_{ij}) - V_t^-(r_{ij})) \hat{P}_{23} \right\} \Phi = E\Phi$$

Critical coupling constants

Let g — is a coupling constant, that characterizes intensity of attractive potential. Than $g_{cr}^{(2)}$, $g_{cr}^{(3)}$ and $g_{cr}^{(4)}$, are critical coupling constants where bound state appears in 2-, 3- and 4-fermion system respectively.

If $k = g_{cr}^{(3)}/g_{cr}^{(2)} > 1$, than 3n may exist only in the region, where 2n is bound.

If $k = g_{cr}^{(3)}/g_{cr}^{(2)} < 1$, than 3n may be bound, while 2n is unbound.



Criteria to choose the potential

If we want to describe bound tetraneutron, we must choose the potential, that can satisfy these conditions:

- ▶ The ratio $k^{(4)} = g_{cr}^{(4)}/g_{cr}^{(2)}$ must be lower than 1 (or $k^{(3)} = g_{cr}^{(3)}/g_{cr}^{(2)}$ for trineutron)
- ▶ Low-energy neutron-neutron scattering parameters:
scattering length ($a_{s(nn)} = -18.9$ fm),
effective radius ($r_{0s(nn)} = 2.75$ fm).
- ▶ Binding energies of light nuclei (3H , 4He ...)
- ▶ ... other experimental data

If no potential can satisfy all these conditions, than bound tri- or tetraneutron can not be described with pairwise potentials.

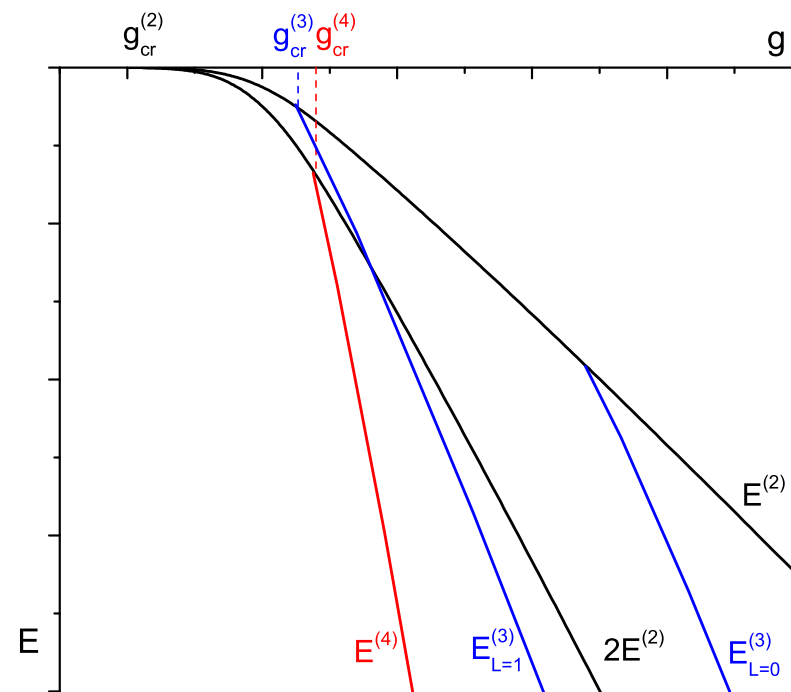
Standard potentials (spinless case)

- ▶ Simplest potential (Gaussian well)

$$V_{ij}(r) = -g e^{-r^2}$$

$$k^{(3)} = 1.43$$

$$k^{(4)} = 1.46$$



This behavior of energy levels is similar for all standard nuclear potentials

Standard potentials (spinless case)

- ▶ Simplest potential (Gaussian well)

$$V_{ij}(r) = -g e^{-r^2}$$

$$k^{(3)} = 1.43$$

$$k^{(4)} = 1.46$$

- ▶ More realistic nuclear potentials with short range repulsion

- Volkov potential:

$$V_{ij}(r) = g (144.86 e^{-(r/0.82)^2} - 83.34 e^{-(r/1.6)^2})$$

$$k^{(3)} = 1.39$$

$$k^{(4)} = 1.44$$

- Best achieved value of k :

$$V_{ij}(r) = g (1.5 e^{-(r/0.9)^2} - e^{-r^2})$$

$$k^{(3)} = 1.278$$

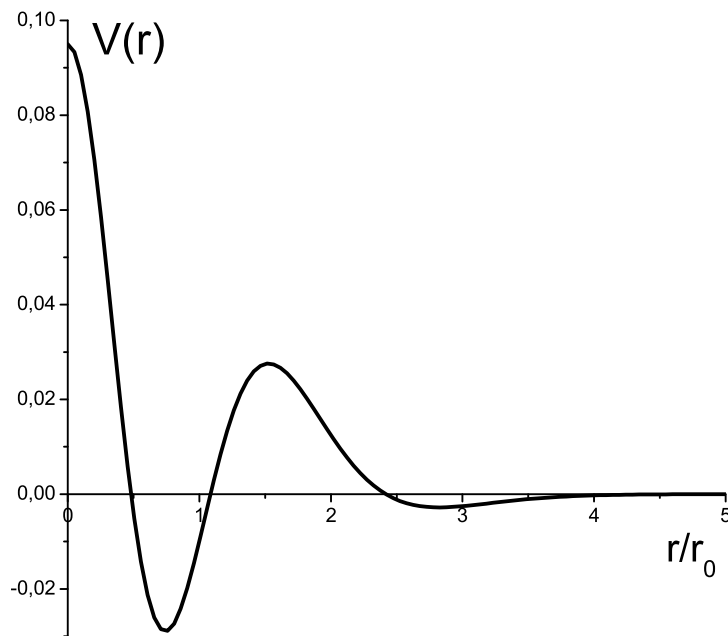
$$k^{(4)} = 1.27$$

There are no possibilities to bind the 3_n or 4_n system with standard or near-standard nuclear pairwise potentials.

Specially constructed potential (spinless case)

- Potential with two attractive regimes

$$V_{ij}(r) = g \left(0.43 e^{-(r/0.6)^2} - e^{-r^2} + 1.085 e^{-(r/1.3)^2} - 0.42 e^{-(r/1.5)^2} \right)$$



This potential gives

$$k^{(3)} = 0.967 < 1 \text{ and}$$

$$k^{(4)} = 0.949 < 1.$$

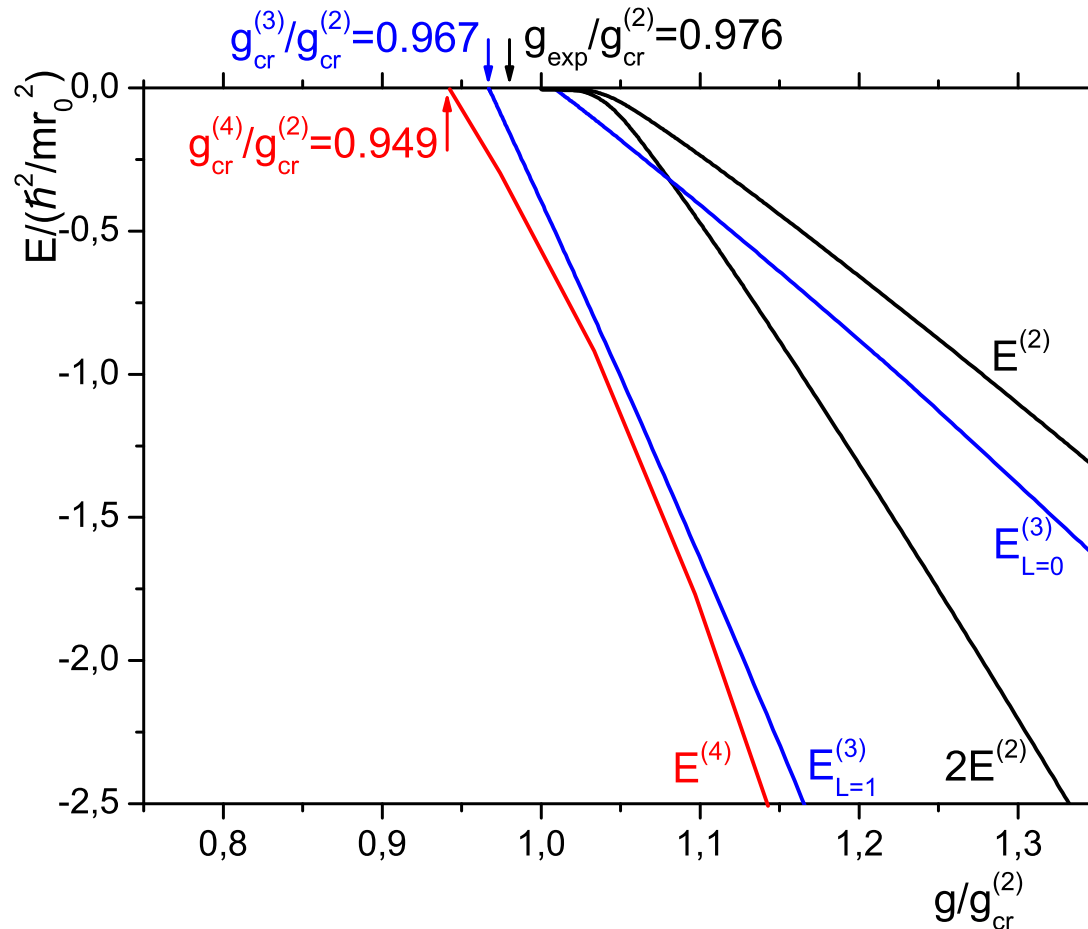
Low-energy n - n parameters are reproduced by proper change to dimensional quantities

$$\left(r \rightarrow \frac{r}{r_0}, V \rightarrow V \frac{2m r_0^2}{\hbar^2}, \text{ where } r_0 = 0.488515 \right)$$

and putting $g = 322.4 = 0.976 g_{cr}^{(2)}$.

This is the only type of n - n potential, that meets criteria above, and hence the only possibility to describe bound 4_n with pairwise potentials.

Behavior of energy levels



Energies of the 4-fermion, 3-fermion (with $L=0$) and 2-fermion systems versus the coupling constant of potential.

► On the graph g_{exper} is a value of coupling constant where low-energy n - n parameters are reproduced. At this point 3- and 4-fermion systems are bound.

► Energy of the 4- and 3-fermion systems are almost parallel to the corresponding thresholds (clustering) in a wide range of coupling constants.

► Slope of the $E^{(4)}$ curve is lower than of $E_{L=1}^{(3)}$ near critical values. It results in that there is a point where these curves are closest to each other.

Spin-dependent potential

Let's choose potential described above as $V_s^+(r)$.

Then putting $V_t^-(r) \equiv 0$ gives no bound 3n and 4n , while we know, that it must be repulsive.

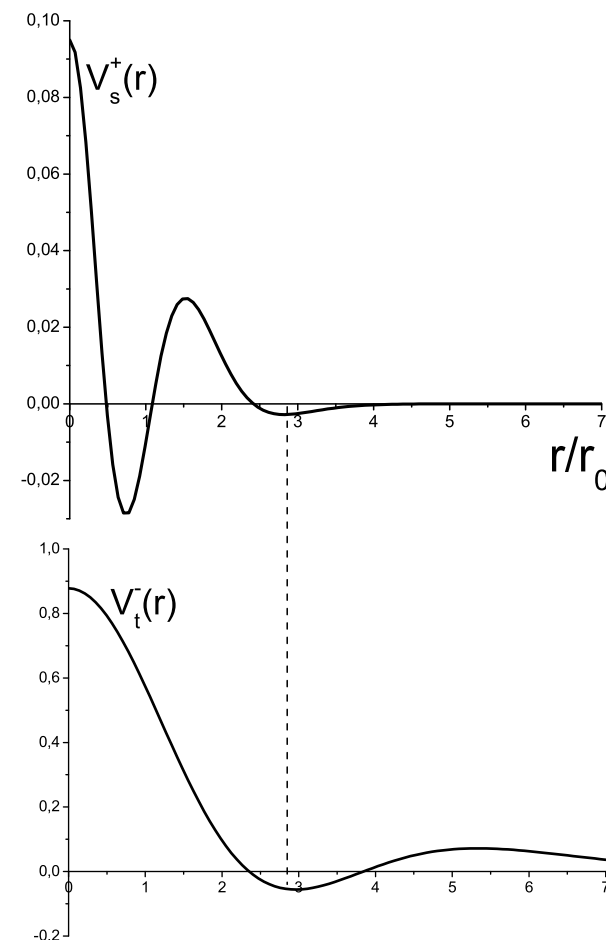
The only way out is to choose repulsive triplet potential with small attractive well, correlated with outer well in singlet.

Full spin-dependent potential:

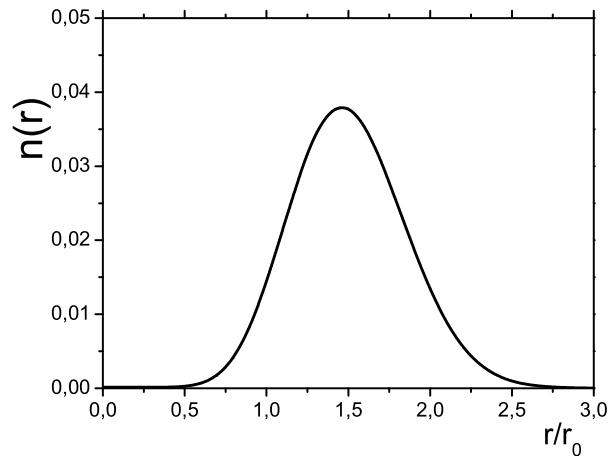
$$V_s^+(r) = g \left(0.43 e^{-(r/0.6)^2} - e^{-r^2} + 1.085 e^{-(r/1.3)^2} - 0.42 e^{-(r/1.5)^2} \right)$$

$$V_t^-(r) = 30.968 e^{-(r/2)^2} - 32.676 e^{-(r/3)^2} + 14.0 e^{-(r/4)^2}$$

This potential gives bound 4n with energy about 0.5 MeV and no bound 3n ($k^{(3)} = 4.27$). There is no possibility to bind 3n with "normal" triplet potential.

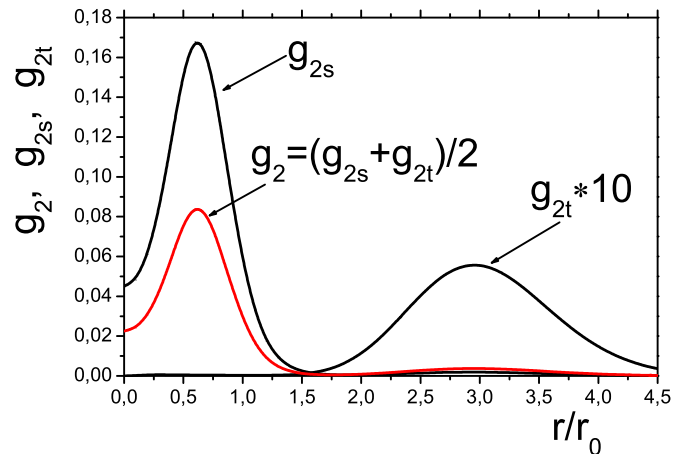


Structure functions of hypothetical 4n



Profile of the density distribution in hypothetical 4n system

- ▶ Anomalously small size of 4-neutron system ($\langle r^2 \rangle^{1/2} = 0.83$ fm)
- ▶ A "bubble" structure of the system.



Profiles of the singlet, triplet, and total pair correlation functions of the hypothetical 4n system.

- ▶ Singlet correlation function is much larger than triplet
- ▶ Singlet correlation function has high peak correlated with inner well in potential
- ▶ Triplet correlation function is nonzero only in the region near attractive well it triplet potential. It means that triplet potential acts as an effective attraction.

Conclusion

- ▶ We develop precise variational schemes and perform reliable calculations of 3- and 4-neutron system. We show, that bound tetraneutron is impossible with standard pairwise potentials.
- ▶ We discovered a wide class of n - n potentials with two attraction regimes that possess following properties:
 - negative value of 4-fermion energy at g_{exper} (bound tetraneutron with $B \simeq 0.5$ MeV)
 - small (~ 2 MeV) **underbinding of 3H and 4He** . Exact energies of such systems can be obtained by slight variation of potentials
 - anomalously small size ($\langle r^2 \rangle^{1/2} = 0.83$ fm) and "bubble" structure of the 4-neutron system
 - high maximum in the singlet scattering phase shift $\delta_s \approx 160^\circ$ at the energies of neutrons near 100 – 150 MeV.
- ▶ Trineutron system can not be bound with any reasonable spin-dependent interaction.
- ▶ More realistic results can not be obtained using pairwise potentials and better description may be achieved only using more complicated types of interactions.

THANK YOU