

COMPLETE POLARIZATION EXPERIMENTS

for $NN \rightarrow N\Delta$ and $p+d \rightarrow (pp)(^1S_0) + \Delta^0$

IN COLLINEAR KINEMATICS

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Content:

- Triple spin correlations for binary reaction $\vec{1} + \vec{2} \rightarrow \vec{3} + 4$
- $\vec{N}\vec{N} \rightarrow N\vec{\Delta}$:
four spin amplitudes $\rightarrow$ observables
- $\vec{p} + \vec{d} \rightarrow pp(^1S_0) + \vec{\Delta}$:
 - three spin amplitudes & observables $\rightarrow$ C.P.E.
 - ★ transversal,
 - ★ longitudinal
- Relation between $\vec{p} + \vec{d} \rightarrow pp(^1S_0) + \vec{\Delta}$ and $NN \rightarrow N\Delta$ within the **IA**

- Δ -isobar is very important for:

$NN \rightarrow NN$ due to coupling to $NN \rightarrow N\Delta$ & $NN \rightarrow \Delta\Delta$

$NN \rightarrow NN\pi$ systems (even near the threshold and above)

3NF forces ($pd \rightarrow pd$ Sagara puzzle; $\frac{d\sigma}{d\Omega}$, 60 - 200 MeV, A_y)

3NF forces in medium

Δ -mechanism in pd -backward elastic scattering, $T > 400\text{MeV}$ T_{20} problem

$\Delta\sigma_L$ & $\Delta\sigma_T$ in pn -, pp - interaction below 1 GeV

Amplitude analysis of the exp.data on $NN \rightarrow N\Delta$ at 500-800 MeV

$(d\sigma/d\Omega, A_{N0}, A_{L0}, A_{S0}, A_{NN}, A_{LL}, P_y\rho_{33})$:

A.B. Wicklund et al., Phys. Rev. D **34** (1986) 19.

R.L. Shypit et. al., Phys. Rev. Lett. **60** (1988) 901; Phys. Rev. C **40** (1989) 2203.

D.V. Bugg, Proceedings of the 3rd International Symposium on πN and NN Physics, LNPI, Gatchina (1989) p.57.

V.V. Anisovich, A.V. Sarantsev, and D.V. Bugg, Nucl. Phys. **A537** (1992) 501.

- Complete polarization experiment for $NN \rightarrow NN$ (5 amplitudes) $\Rightarrow$ SAID
- 16 amplitudes for $NN \rightarrow N\Delta$ in general case
- Collinear kinematics (only one direction $\mathbf{k}$) $\Rightarrow$ simplification of spin dependence
Polarizations with axial symmetry is not a limiting case of the general formalism

Examples:

- $dp \rightarrow pd$ in collinear kinematics

V.P. Ladygin, N.B. Ladygina, J.Phys.G: Part.Nucl. **23** (1997) 847

M.P. Rekalo, N.M. Piskunov, and I.M. Sitnik, Few-Bod.Syst. **23** (1998) 187.

Review of the near threshold formalism:

E. Tomasi-Gustaffson and M.P. Rekalo, Fiz. Elem. Chat. At.Yadr. **33** (2002) 436
(Phys. Part.Nucl. **33** (2002) 220).

$$\underline{\vec{1} + \vec{2} \rightarrow \vec{3} + 4}$$

Beam direction $\mathbf{k}||OZ$

$$T_{\mu_1 \mu_2}^{\mu_3 \mu_4} = \langle \mu_1, \mu_2 | F | \mu_3 \mu_4 \rangle = \sum_{S_i M_i S_f M_f L m} (j_1 \mu_1 j_2 \mu_2 | S_i M_i) (j_3 \mu_3 j_4 \mu_4 | S_f M_f) \times \\ \times (S_i M_i L m | S_f M_f) Y_{Lm}(\hat{\mathbf{k}}) a_{S_f}^{L S_i}. \quad (1)$$

$$(-1)^L = \pi_1 \pi_2 \pi_3 \pi_4,$$

Triple spin correlation

$$K_{J_1 M_1, J_2 M_2}^{J_3 M_3} = \frac{Tr \{ T_{J_3 M_3}(3) F T_{J_1 M_1}(1) T_{J_2 M_2}(2) F^+ \}}{Tr F F^+}, \quad (2)$$

Rank $J_i = 0, 1, \dots, 2j_i$

$$M_k = -J_k, -J_k + 1, \cdots, J_k$$

S.M.Bilenky, L.I.Lapidus, L.D.Puzikov, R.M. Ryndin, Nucl.Phys. **7** (1958) 646.

$$\begin{aligned}
& \textcolor{red}{K}_{J_1 M_1, J_2 M_2}^{J_3 M_3} \text{Tr} F F^+ = \frac{1}{4\pi} \sqrt{(2J_1 + 1)(2J_2 + 1)(2J_3 + 1)} \\
& \sum_{S S' J J' L L' J_0 J'_0} (2J + 1)(2J' + 1) \sqrt{(2L + 1)(2L' + 1)} \times \\
& \sqrt{(2S + 1)(2S' + 1)(2J_0 + 1)} (-1)^{j_3 + j_4 + J + L + S' - S} \times \\
& \times (J_0 - M_3 J_3 M_3 | J'_0 0) (L' 0 L 0 | J'_0 0) (J_1 M_1 J_2 M_2 | J_0 - M_3) \times \\
& \left\{ \begin{matrix} j_3 & j_4 & J \\ J' & J_3 & j_3 \end{matrix} \right\} \left\{ \begin{matrix} S & j_1 & j_2 \\ S' & j_1 & j_2 \\ J_0 & J_1 & J_2 \end{matrix} \right\} \left\{ \begin{matrix} S & J & L \\ S' & J' & L' \\ J_0 & J_1 & J'_0 \end{matrix} \right\} \textcolor{red}{a}_J^{L S} (\textcolor{red}{a}_{J'}^{L' S'})^*. \tag{3}
\end{aligned}$$

$$\begin{aligned}
& \textcolor{blue}{L} + \textcolor{blue}{L}' + \textcolor{blue}{J}'_0 \text{ is even, } \textcolor{blue}{J}'_0 \text{ is even} \\
& \textcolor{red}{M}_1 + \textcolor{red}{M}_2 + \textcolor{red}{M}_3 = 0. \tag{4}
\end{aligned}$$

Symmetry:

$$\textcolor{blue}{K}_{J_1 - M_1, J_2 - M_2}^{J_3 - M_3} = (-1)^{J_1 + J_2 + J_3} \textcolor{blue}{K}_{J_1 M_1, J_2 M_2}^{J_3 M_3} \tag{5}$$

• The reaction $NN \rightarrow \Delta N$

Parity and angular momentum conservation $\rightarrow a_{S_f}^{L S_i}$:

$$B_1 = a_2^{20}, \underbrace{B_2 = a_2^{21}, B_3 = a_1^{21}, B_4 = a_1^{01}}_{\text{(spin-triplet)}}$$

$$|B_1|^2 = \left\{ \frac{1}{5}(K_{10,10}^{20} - K_{00,00}^{20}) + \frac{2}{5}K_{11,1-1}^{00} \right\} \Sigma, \quad (6)$$

$$|B_2|^2 = \frac{1}{5} \{ 1 + K_{00,00}^{20} - \sqrt{6}K_{11,11}^{2-2} - 3K_{10,10}^{20} \} \Sigma, \quad (7)$$

$$|\sqrt{2}B_3 - B_4|^2 = \{ K_{10,10}^{20} - K_{00,00}^{20} - 2K_{11,1-1}^{00} \} \Sigma, \quad (8)$$

$$|B_3 + \sqrt{2}B_4|^2 = \frac{1}{2} \left\{ 1 + 3K_{00,00}^{20} - K_{10,10}^{20} + \sqrt{6}K_{11,11}^{2-2} \right\} \Sigma, \quad (9)$$

where

$$\Sigma = 4\pi T r F F^+ = \frac{16\pi d\sigma_0}{\Phi}. \quad (10)$$

$$Re B_2 (B_3 + \sqrt{2}B_4)^* = \frac{\Sigma}{2\sqrt{10}} \{ \sqrt{6}(K_{10,10}^{20} + K_{00,00}^{20}) - 2K_{11,11}^{2-2} \}, \quad (11)$$

$$Re B_1 (\sqrt{2}B_3 - B_4)^* = \Sigma \frac{1}{30} \{ \sqrt{5} + 20K_{10,00}^{30} - 2\sqrt{5}K_{00,00}^{20} + 3K_{10,10}^{20} \}. \quad (12)$$

For **T-odd observables**($J_1 + J_2 + J_3$ is odd):

$$K_{11,11}^{3-2} \Sigma = \frac{i5\sqrt{2}}{2} \text{Im} B_2 (B_3 + \sqrt{2}B_4)^*, \quad (13)$$

$$K_{1-1,11}^{30} \Sigma = \frac{i3}{2} \text{Im} B_1 (\sqrt{2}B_3 - B_4)^* \quad (14)$$

Real and imaginary parts of the products $B_1(\sqrt{2}B_3 - B_4)^*$ and $B_2(B_3 + \sqrt{2}B_4)^* \Rightarrow$ the relative phases $\phi_{B_1} - \phi_{\sqrt{2}B_3 - B_4}$ and $\phi_{B_2} - \phi_{B_3 + \sqrt{2}B_4}$.
The last relative phase, for example, $\kappa = \phi_{B_2} - \phi_{B_3 + \sqrt{2}B_4}$, can be found from $K_{11,10}^{3-1}$ and $K_{11,00}^{3-1} \Rightarrow \sin \kappa$ and $\cos \kappa$.

Complete polarization experiment:

$$d\sigma_0, \underbrace{K_{11,1-1}^{00}, K_{10,10}^{20}, K_{11,11}^{2-2}, K_{00,00}^{20}, K_{10,00}^{30}, K_{11,00}^{31}}_{\text{T-even,}} \underbrace{K_{11,11}^{3-2}, K_{1-1,11}^{30} \text{ and } K_{11,10}^{3-1}}_{\text{T-odd}}.$$

The reaction $pd \rightarrow \Delta^0(pp)(^1S_0)$

$$A_1 = a_{\frac{3}{2}}^{2\frac{1}{2}}, A_2 = a_{\frac{3}{2}}^{0\frac{1}{2}}, A_3 = a_{\frac{3}{2}}^{2\frac{1}{2}}$$

★ Transversally polarized beam and target

Six observables, $d\sigma_0$, $K_{00,20}^{00}$, $K_{1-1,11}^{00}$, $K_{00,00}^{20}$, $K_{00,21}^{2-1}$, and $K_{00,2-2}^{22}$, $\Rightarrow$
 $\Rightarrow |A_1|^2$, $|A_2|^2$, $|A_3|^2$ and cosines of $\phi_1 - \phi_2$ and $\phi_1 - \phi_3$:

$$\begin{aligned} |A_1|^2 &= I_0 + \frac{2E + F}{3}, \\ |A_2|^2 &= \frac{1}{2}C - \frac{1}{3}E - \frac{1}{6}F, \\ |A_3|^2 &= -\frac{1}{2}C - \frac{1}{3}E - \frac{1}{6}F, \\ \text{Re}A_1A_2^* &= \frac{1}{2}D + \frac{E - F}{6}, \\ \text{Re}A_1A_3^* &= \frac{1}{2}D - \frac{E - F}{6}, \end{aligned} \tag{15}$$

where

$$I_0 = |A_1|^2 + |A_2|^2 + |A_3|^2 = \pi \text{Tr} F F^+ = \frac{6\pi d\sigma_0}{\Phi}, \tag{16}$$

$$\begin{aligned}
E &= \{4K_{1-1,11}^{00} - \frac{4}{\sqrt{3}}K_{00,20}^{00} - \frac{2}{3}\}I_0, \\
F &= \{\frac{12}{\sqrt{3}}K_{00,20}^{00} - 2\sqrt{6}K_{00,00}^{20} - 1\}I_0, \\
C &= \{\frac{4}{\sqrt{3}}K_{00,2-2}^{22} - \frac{8}{\sqrt{3}}K_{00,21}^{2-1}\}I_0, \\
D &= -\frac{4}{\sqrt{3}}\{K_{00,2-2}^{22} + K_{00,21}^{2-1}\}I_0.
\end{aligned} \tag{17}$$

In order to find the sines of two phases one should measure the three T-odd observables $K_{00,11}^{2-1}$, $K_{11,00}^{2-1}$, and $K_{11,2-1}^{00}$.

$$\begin{aligned}
iImA_1A_2^* &= -2\{K_{11,2-1}^{00} - \frac{1}{\sqrt{3}}K_{00,11}^{2-1} + \frac{1}{\sqrt{2}}K_{11,00}^{2-1}\}I_0, \\
iImA_1A_3^* &= 2\{K_{11,2-1}^{00} + \frac{1}{\sqrt{3}}K_{00,11}^{2-1} - \frac{1}{\sqrt{2}}K_{11,00}^{2-1}\}I_0,
\end{aligned} \tag{18}$$

For $J_1 + J_2 + J_3$ being odd, the value of $K_{J_1M_1,J_2M_2}^{J_3M_3}$ is purely imaginary.

Note: only the signs of $ImA_1A_2^*$ and $ImA_1A_3^*$ are required.

Nine observables $\Rightarrow$ complete polarization experiment: $d\sigma_0$, $K_{00,20}^{00}$, $K_{1-1,11}^{00}$, $K_{00,00}^{20}$, $K_{00,21}^{2-1}$, $K_{00,2-2}^{22}$, $\underbrace{K_{00,11}^{2-1}, K_{11,00}^{2-1}, K_{11,2-1}^{00}}_{\text{T-odd}}$.

★ Transversally & longitudinally polarized beam and target in $pd \rightarrow \Delta^0(pp)(^1S_0)$

Using I_0 and $\underbrace{K_{10,10}^{00}}_{\text{longitudinal}} \underbrace{K_{1-1,11}^{00}, K_{00,2-1}^{21}, K_{00,2-2}^{22}}_{\text{transversal}} \Rightarrow |A_1|^2, |A_2|^2, \text{ and } |A_3|^2.$

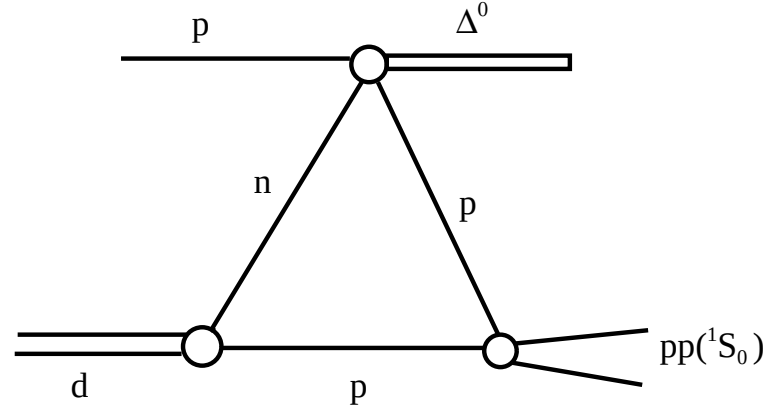
In addition, $K_{00,00}^{20} \Rightarrow \text{Re } A_1 A_2^*, \text{Re } A_1 A_3^* \text{ and } \text{Re } A_2 A_3^*.$

Finally, $K_{00,11}^{2-1}$ and $K_{11,00}^{2-1}$ give

$$i\text{Im} A_2 A_3^* = \left\{ \frac{4}{\sqrt{3}} K_{00,11}^{2-1} + \frac{2}{\sqrt{2}} K_{11,00}^{2-1} \right\} I_0 \quad (19)$$

$\phi_{23} = \phi_{13} + \phi_{21}$, where $\phi_{ij} = \phi_i - \phi_j$

The above **eight observables** $\Rightarrow A_1, A_2$ and A_3



$$B_i = a_J^{LS}(pn \rightarrow \Delta^0 p) \rightarrow M_{\mu_0, \lambda}^{\mu \Delta}(pd \rightarrow \Delta^0(pp)_{1S_0})$$

$$\begin{aligned} A_1 &= \frac{\sqrt{6\pi}}{3}(M_{+,0}^{+1/2} - \sqrt{2}M_{+,-}^{-1/2}), \\ A_2 &= \frac{\sqrt{6\pi}}{3}\left(\sqrt{\frac{3}{2}}M_{+,+}^{+3/2} + M_{+,0}^{+1/2} + \frac{1}{\sqrt{2}}M_{+,-}^{-1/2}\right), \\ A_3 &= \frac{\sqrt{6\pi}}{3}\left(\sqrt{\frac{3}{2}}M_{+,+}^{+3/2} - M_{+,0}^{+1/2} - \frac{1}{\sqrt{2}}M_{+,-}^{-1/2}\right), \end{aligned} \tag{20}$$

A_1, A_2, A_3 are functions of B_1, B_2, B_3, B_4 and transition form factors

Inversion problem:

$$\begin{aligned}
B_4 &= -\frac{16\pi}{\sqrt{15}}\{(5S_2 + \sqrt{10}S_0)A_2 + 5S_2(A_3 - \sqrt{3}A_1)\}/Y, \\
\sqrt{5}B_2 - \sqrt{3}B_3 &= \frac{32\pi}{\sqrt{2}}\{\sqrt{5}S_2(A_1 + A_2) + (\sqrt{2}S_0 + \sqrt{5}S_2)A_3\}/Y, \\
\sqrt{3}B_1 - \sqrt{2}B_2 &= \frac{16\pi}{\sqrt{5}}\{(2\sqrt{2}S_0 + \sqrt{5}S_2)A_1 - 3\sqrt{5}S_2A_2 - (\sqrt{2}S_0 - \sqrt{5}S_2)A_3\}/Y,
\end{aligned} \tag{21}$$

where

$$Y = 2S_0^2 + \sqrt{10}S_0S_2 - 10S_2^2. \tag{22}$$

Transition form factor ${}^3S_1 - {}^3D_1 \rightarrow {}^1S_0$

$$S_l(Q/2) = \int_0^\infty dr r^2 j_l(Qr/2) u_l(r) \psi_k^{(-)*}(r), \tag{23}$$

where u_0 and u_2 are the S- and D-components of the deuteron $\psi_k^{(-)}(r)$ is the NN scattering wave function in the 1S_0 state

CONCLUSION

D.V. Bugg and C. Wilkin, Nucl. Phys. A467 (1987) 575.

Charge-exchange $dp \rightarrow (pp)_{(1S_0)}n$ and $pn \rightarrow np$ (SATURNE, COSY, JINR)

- ★ Complete polarizations experiment for $NN \rightarrow N\Delta$ (4 amplitudes, 10 observables)
- ★ Complete polarizations experiment for $pd \rightarrow (pp)_{(1S_0)}\Delta^0$ (3 amplitudes, 8-9 observables)
- ★ Three combinations from four amplitudes of the $NN \rightarrow N\Delta$ are related to three amplitudes of the $pd \rightarrow (pp)_{(1S_0)}\Delta^0$

”... The whole business looks extremely complicated — but that is the nature of the problem.”

$$\left(\frac{3}{2} \neq \frac{1}{2}\right)$$

C. Wilkin (private communication)